

The Hipparcos observations and the mass of sub-stellar objects^{*}

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Abstract. The Hipparcos Intermediate Astrometric Data have been used lately to estimate the inclination of the orbital plane of candidate extrasolar planets. Whereas most of these investigations derive almost face-on orbits, we show that the astrometric data are seldom precise enough to undertake such studies and that the “face-on” result might be just a spurious effect of the method.

Key words. methods: data analysis – astrometry – stars: planetary systems

1. Introduction

Even before the release of the Hipparcos and Tycho Catalogues (ESA 1997), their interest regarding the extrasolar planets was already pointed out (Perryman et al. 1997). For the past two years, some investigators have re-processed the Hipparcos Intermediate Astrometric Data (IAD) of different classes of objects. Though initially limited to binaries (Söderhjelm 1999; Halbwegs et al. 2000; Pourbaix & Jorissen 2000), the planet hunters have also got interested in such re-processings (Mazeh et al. 1999; Gatewood et al. 2000; Zucker & Mazeh 2000). A recent statistical study based on the IAD concluded that a significant fraction of the orbits of extrasolar planets are seen almost face-on (Han et al. 2001), thus pushing up their mass estimates. According to that study, a substantial number of these objects would no longer be planets.

We have re-analyzed the IAD of 46 stars with planetary candidates and found that the orbital model, with the spectroscopic parameters assumed, seldom improves the astrometric fit significantly. We show that in most cases the small inclinations found by Han et al. (2001) are just an artifact of the fitting procedure. Such a caveat is indeed related to the size of the orbit with respect to the precision of the astrometric measurements. Similar results, although wrong, would therefore be derived with classes of instruments of the same given precision (e.g. MAP, Gatewood 1987).

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2. Astrometric model and the choice of the parameters

Regardless of its physical nature, the motion of any member of a binary system with respect to the common center of mass is described by the well-known relations (Binnendijk 1960)

$$x = A(\cos E - e) + F\sqrt{1 - e^2} \sin E, \quad (1)$$

$$y = B(\cos E - e) + G\sqrt{1 - e^2} \sin E \quad (2)$$

where E is the eccentric anomaly, solution of Kepler’s equation for given e , P , and T . Seven orbital parameters are thus required: A , B , F , G , e , P , and T . The first four, the Thiele-Innes constants, are linked to the Campbell elements (a_a , i , ω_1 , and Ω) through

$$A = a_a(\cos \omega_1 \cos \Omega - \sin \omega_1 \sin \Omega \cos i), \quad (3)$$

$$B = a_a(\cos \omega_1 \sin \Omega + \sin \omega_1 \cos \Omega \cos i), \quad (4)$$

$$F = a_a(-\sin \omega_1 \cos \Omega - \cos \omega_1 \sin \Omega \cos i), \quad (5)$$

$$G = a_a(-\sin \omega_1 \sin \Omega + \cos \omega_1 \cos \Omega \cos i). \quad (6)$$

There is a priori no mathematical reason to prefer the Campbell elements over the Thiele-Innes ones. However, the spectroscopic orbit yields, among other quantities, ω_1 and the product $a_1 \sin i$ (a_1 is here reckoned in linear units whereas a_a is its angular counterpart). Therefore, it is convenient to use Campbell’s elements even if global optimization techniques are then required to minimize χ^2 (Pourbaix & Jorissen 2000) because the model is highly non-linear in terms of the remaining elements.

Most of the sub-stellar objects have been detected through radial velocity surveys and, therefore, the spectroscopic orbit of the primary is known. Coupled to the

parallax (ϖ), two parameters thus remain unknown, namely i and Ω . Indeed, a_a can be written as

$$a_a = 3.35729138 \cdot 10^{-5} K_1 P \sqrt{1 - e^2} \varpi / \sin i, \quad (7)$$

where K_1 is the amplitude of the radial velocity curve reckoned in m s^{-1} , P is the period in yr, ϖ in mas and e is the eccentricity.

The IAD were made available to allow some further processing of the Hipparcos observations, especially those of binaries. Regardless of the model used to derived the final Hipparcos results, the IAD are residuals with respect to the single star model. When the binary nature of a star was discovered after the release of the Catalogues, the IAD can therefore be re-processed with an orbital model (Halbwachs et al. 2000; Pourbaix & Jorissen 2000). Although they are applied to 1-dimensional data, the equations of the orbital perturbation are essentially those of the visual absolute orbit. They thus supply with the inclination i .

It is very tempting to re-process the IAD of all the stars around which planets were lately discovered. Indeed, the motion of these stars should reveal tiny astrometric perturbations which can now be explained by the planet orbiting them. Fitting the IAD should supply with the inclination and therefore with the mass of the companion. For instance, that technique was applied to HD 10697 (HIP 8159) to obtain the mass of a brown dwarf (Zucker & Mazeh 2000). The brown dwarf nature of the companion was derived from the rather low value of the inclination, thus leading to a quite large mass for the secondary.

3. Inclination from a tiny orbital wobble

Why does fitting i and Ω yield almost face-on orbits whereas, from a naive reasoning, one would expect preferentially edge-on situations for spectroscopically detected companions? Halbwachs et al. (2000) barely mention that feature. Indeed, in order to validate their approach, they process the IAD of known single stars with the orbital model (with ω_1 , e , P , and T fixed) and investigate the behavior of a_a . Instead of $a_a = 0$ (the real value for single stars), they derive its Rayleigh distribution and conclude that a_a is slightly overestimated (when the same method is applied to genuine binaries).

They also noticed that the estimated a_a is directly related to the residuals of the coordinates of the star, i.e. about 1 mas for Hipparcos. Fitting an orbit to the IAD when $a_a \sin i \ll 1$ mas always leads to small inclinations: the smaller the product, the smaller the inclination. This is only due to the relative precision of the instrument with respect to the magnitude of the orbital wobble. The criterion depends on $a_a \sin i$, not on K_1 . Indeed, long-period astrometric binaries also characterized by small K_1 would be detected with a much longer Hipparcos-type mission because, *mutatis mutandis*, $a_a \sin i$ would increase with P . The orbital periods of known sub-stellar candidates are simply too small for $a_a \sin i$ to be large enough.

What about the “precision” of such inclinations? Let us have a look at the derivative of the Thiele-Innes constant A with respect to i (the reasoning is the same for B , F , and G). These derivatives are indeed used to build the Fisher matrix whose inverse is the covariance matrix of the fitted parameters (Bevington & Robinson 1992).

$$\frac{\partial A}{\partial i} = \frac{\partial a_a}{\partial i} (\cos \omega_1 \cos \Omega - \sin \omega_1 \sin \Omega \cos i) + a_a \sin \omega_1 \sin \Omega \sin i. \quad (8)$$

If a_a is an independent parameter in the astrometric solution, the first term disappears and the second vanishes with i . In that case, the smaller the inclination, the larger its uncertainty. On the other hand, when Eq. (7) is adopted as a constraint,

$$\frac{\partial a_a}{\partial i} = -a_a \cot i. \quad (9)$$

Hence, the smaller the inclination, the larger the derivatives, the smaller the standard deviation of the inclination, regardless of the precision of the astrometric data.

4. Do these orbits improve the fit?

Fitting the IAD of HD 209458 (HIP 108859) using the spectroscopic orbit by Mazeh et al. (2000) yields $i \approx 0.02^\circ$ even if we know that transits do occur (Charbonneau et al. 2000), implying $i \sim 90^\circ$. We need a way to rule out such solutions.

Table 1 lists extrasolar planet candidates with their $a_a \sin i$ based on their Hipparcos parallax and their most recent spectroscopic orbit. For all but seven, the value of $a_a \sin i$ constrains the inclination to be $< 10^\circ$ (or $> 170^\circ$) if one assumes that Hipparcos noticed the orbital motion. That latter assumption is where the weakness of the reasoning lies. An F-test can be performed to see whether the two additional parameters improve the fit of the IAD with respect to the single star model. In the absence of an astrometric wobble, the quantity

$$\hat{F} = \frac{N - 7}{2} \frac{\chi_S^2 - \chi_C^2}{\chi_C^2} \quad (10)$$

follows the F-distribution with $(2, N - 7)$ degrees of freedom (Bevington & Robinson 1992). N denotes the number of data points and χ_S^2 and χ_C^2 are the value of the χ^2 with the 5-parameter (single star) model and orbital model respectively.

The column labeled α in Table 1 gives the probability of obtaining an F -value greater or equal to \hat{F} if the null hypothesis H_0 : “no orbital wobble present in the IAD” holds. That hypothesis is rejected for all but four stars at a 5% level. Even if two rejections are expected by chance in this sample, HD 38529 (HIP 27253), HD 83443 (HIP 47202), ρ CrB (HIP 78459), and HD 195019 (HIP 100970) do deserve some further investigations. However, a small α only means that an orbital model improves the fit of the

Table 1. List of planet candidates and their $a_a \sin i$ based on the spectroscopic orbit and the Hipparcos parallaxes. i is the inclination derived from the IAD when $a_a \sin i$ is used as a constraint. α is the probability of obtaining an F -value greater or equal to \hat{F} if there is no orbital wobble present in the IAD. Ref. is the reference for the orbit: 1: Naef et al. (2000); 2: Queloz et al. (2000a); 3: Butler et al. (1999); 4: Vogt et al. (2000); 5: Fischer et al. (2001); 6: Queloz et al. (2000b); 7: Marcy et al. (2000); 8: Kürster et al. (2000); 9: Hatzes et al. (2000); 10: Butler et al. (2000); 11: Butler et al. (1997); 12: Naef et al. (2001); 13: Mayor et al. (2000); 14: Korzennik et al. (2000); 15: Butler & Marcy (1996); 16: Marcy et al. (1999); 17: Marcy & Butler (1996); 18: Udry et al. (2000a); 19: Noyes et al. (1999); 20: Udry et al. (2000b); 21: Cochran et al. (1997); 22: Sivan et al. (2000); 23: Mazeh et al. (2000); 24: Delfosse et al. (1998); 25: Mayor & Queloz (1995)

HIP	HD/ Name	Ref.	$a_a \sin i$ (mas)	i ($^\circ$)	α (%)	HIP	HD/ Name	Ref.	$a_a \sin i$ (mas)	i ($^\circ$)	α (%)
1292	GJ 3021	1	9.86e-02	12.3 \pm 13.6	65	64426	114762	16	1.18e-01	4.3 \pm 3.4	41
5054	6434	2	1.79e-03	0.2 \pm 0.2	51	65721	70 Vir	17	1.72e-01	13.7 \pm 9.5	42
7513b	ν And b	3	2.30e-03	179.7 \pm 0.3	62	67275	τ Boo	11	9.28e-03	0.9 \pm 0.8	55
7513c	ν And c	3	9.43e-02	173.7 \pm 3.8	23	68162	121504	2	6.16e-03	0.3 \pm 0.2	27
7513d	ν And d	3	5.76e-01	28.7 \pm 16.8	20	72339	130322	18	3.83e-03	0.2 \pm 0.2	61
8159	10697	4	3.48e-01	169.2 \pm 4.2	6	74500	134987	4	4.41e-02	2.6 \pm 2.8	55
9683	12661	5	5.55e-02	3.0 \pm 1.9	36	78459	ρ CrB	19	1.35e-02	179.1 \pm 0.5	1
10138	Gl 86	6	5.06e-02	8.5 \pm 11.9	75	79248	14 Her	20	7.75e-01	140.6 \pm 27.1	25
12048	16141	7	2.08e-03	0.1 \pm 0.1	46	89844b	168443b	20	5.95e-02	178.0 \pm 1.8	50
12653	HR 810	8	1.14e-01	7.0 \pm 4.4	20	89844c	168443c	20	1.11e+00	48.7 \pm 42.4	60
14954	19994	2	8.34e-02	4.9 \pm 3.7	40	90485	169830	12	4.60e-02	2.1 \pm 1.1	10
16537	ϵ Eri	9	1.06e+00	174.0 \pm 4.1	44	93746	177830	4	1.86e-02	1.3 \pm 0.7	19
26381	37124	4	1.80e-02	179.5 \pm 0.4	55	96901	16 Cyg B	21	1.18e-01	170.8 \pm 7.2	41
27253	38529	5	1.62e-03	0.1 \pm 0.0	3	97336	187123	4	4.09e-04	180.0 \pm 0.0	17
31246	46375	7	2.95e-04	0.0 \pm 0.0	46	98714	190228	22	1.40e-01	5.1 \pm 2.4	9
33719	52265	10	1.66e-02	178.4 \pm 1.5	59	99711	192263	4	7.62e-03	179.4 \pm 0.5	73
43587	55 Cnc	11	8.38e-03	179.5 \pm 0.3	30	100970	195019	4	1.28e-02	0.3 \pm 0.1	< 1
47007	82943	12	8.58e-02	7.2 \pm 7.3	61	108859	209458	23	5.75e-04	0.0 \pm 0.0	19
47202b	83443b	13	3.50e-04	0.0 \pm 0.0	13	109378	210277	4	6.59e-02	175.3 \pm 3.6	41
47202c	83443c	13	7.78e-04	0.0 \pm 0.0	3	113020	Gl 876	24	2.82e-01	172.2 \pm 7.2	49
50786	89744	14	1.06e-01	175.4 \pm 2.9	35	113357	51 Peg	25	1.49e-03	0.1 \pm 0.2	65
52409	92788	2	1.28e-01	4.4 \pm 2.8	22	113421	217107	4	4.58e-03	179.6 \pm 0.4	54
53721	47 UMa	15	3.24e-01	45.4 \pm 72.0	86	116906	222582	4	1.58e-01	5.4 \pm 2.9	10

IAD, with no guaranty that the adopted parameters indeed yield the best possible fit. So the obtained inclination might be unreliable even when $\alpha \approx 0$.

The F-test rejects the astrometric orbital solutions for ν And (HIP 7513) and for HD 10697 as well as the face-on solution we would obtain for HD 209458. It is worth pointing out that the value of α depends on the adopted parameters. Therefore, an alternative orbit might still substantially improve the fit even if the value of α listed in Table 1 is high.

One should also mention that $a_a \sin i$ large with respect to the precision of the instrument is not a sufficient condition for a reliable astrometric solution. For instance, Pourbaix & Jorissen (2000) fitted the IAD of 81 single-lined spectroscopic binaries and obtained reliable results for only 24 of them. For 21 out of these 24, $a_a \sin i$ ranges from 0.6 to 13.8 mas whereas only 20 among the 57 others are characterized by a $a_a \sin i$ much smaller than 1 mas. However, 27% of the stars rejected although they fulfill the 1 mas criterion have periods exceeding 11 years. So no orbital solution could be derived for them due to the poor orbit coverage during the Hipparcos mission.

In the present sample, HD 168443c (HIP 89844) and ϵ Eri (HIP 16537) both have $a_a \sin i > 1$ mas but fail to get a

significant astrometric solution probably because of their relatively long periods (4.57 and 6.86 years, respectively).

5. Conclusion

Fitting (i, Ω) to the IAD when the spectroscopically constrained $a_a \sin i$ is much smaller than the astrometric precision always yields low values of $\sin i$, irrespective of the true inclination. Although one cannot rule out the possibility of almost face-on orbits, very few of these orbits result in a significant improvement of the astrometric fit with respect to the single star model.

Hipparcos was a very successful mission and its files certainly still hold some hidden results. However, if one does not take care, one may be tempted to make these observations tell more than what they actually can. Moreover, the fact that one derives the same result with two instruments belonging to the same class of precision, for instance MAP and Hipparcos, does not always suffice to assess its reliability. Before a $100\mu\text{as}$ -class instrument becomes available, the astrometric techniques will not derive the mass of most of the extrasolar sub-stellar companions known today. The four stars with a formally

significant astrometric orbit ($\alpha < 5\%$ in Table 1) do on the other hand deserve serious further consideration.

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